

Optimization of large CI wave functions in QMC

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 \rightarrow CI wavefunction: linear combination of Slater determinants $\{D_I\}$

$$\Phi(\mathbf{r}_1,\cdots,\mathbf{r}_N)=\sum_{I=1}^{N_{\text{det}}}C_I\,D_I(\mathbf{r}_1,\cdots,\mathbf{r}_N)$$

 \rightarrow Usually, $N_{\rm det} \ll N_{\rm FCL}$. We improve Φ by adding a Jastrow factor:

$$\Psi(\mathbf{r}_1, \dots, \mathbf{r}_N) = \underbrace{\sum_{I=1}^{N_{\text{det}}} C_I D_I(\mathbf{r}_1, \dots, \mathbf{r}_N)}_{\Phi(\mathbf{r}_1, \dots, \mathbf{r}_N)} \exp[J]$$

ightharpoonup goal: optimize $\{C_I\}$ in the presence of Jastrow factor for large $N_{\rm det}$ ($\sim 10^6$)



→ Common formulation of the problem:

$$\widehat{H}\left(e^{J}\Phi\right)=E\left(e^{J}\Phi\right)$$

→ Transcorrelated formalism:

$$\underbrace{\left(e^{-J}\widehat{H}e^{J}\right)}_{\widehat{H}_{\mathsf{TC}}}\Phi = E\Phi$$

- \Rightarrow In practice, we end up with integrals involving $e^{-J}\widehat{H}e^{J}$ instead of $e^{J}\widehat{H}e^{J}$
- \Rightarrow we will optimize the coefficients by solving the TC eigenproblem in the $\{D_I\}$ basis



$$\widehat{H}_{\mathsf{TC}} \Phi = E_{\mathsf{TC}} \Phi \Rightarrow \sum_{K=1}^{N_{\mathsf{det}}} \left\langle D_I \left| \widehat{H}_{\mathsf{TC}} - \widehat{H} + \widehat{H} \left| D_K \right\rangle \right. C_K^{(i)} = E_{\mathsf{TC}} C_I^{(i)}$$

$$\sum_{K \neq I} \mathsf{H}_{IK} C_K^{(i)} + \left(\mathsf{H}_{II} + \left[\frac{1}{C_I^{(i-1)}} \left\langle D_I \left| \widehat{H}_{\mathsf{TC}} - \widehat{H} \right| \Phi^{(i-1)} \right\rangle \right] \right) C_I^{(i)} \approx E_{\mathsf{TC}} C_I^{(i)}$$

$$\underset{\mathsf{dressing elements}}{\underbrace{ \mathsf{dressing elements}}}$$

① we build the diagonal dressing matrix
$$\Delta^{(i-1)}$$
:
$$\Delta_{IK}^{(i-1)} = \begin{cases} \frac{1}{C_I^{(i-1)}} \left[\boxed{\left\langle D_I \middle| \widehat{H}_{\mathsf{TC}} \middle| \Phi^{(i-1)} \right\rangle} - \left\langle D_I \middle| \widehat{H} \middle| \Phi^{(i-1)} \right\rangle \right] & \text{if } I = K \\ 0 & \text{otherwise} \end{cases}$$

- 2 we apply Davidson to extract the new ground state
- ③ iterate until convergence



Advantages of this TC-VMC approach:

- \rightarrow memory scale: $\mathcal{O}(N_{\text{det}})$
- \Rightarrow fast convergence ($\sim 2-3$ iterations) we have applied this approach to optimize $\sim 400\,000$ coefficients with only one iteration.
- \Rightarrow small fluctuations on the involved integrals (compared to $e^{J}\widehat{H}e^{J}$)



Statistical fluctuations: improved estimator

We need to calculate N_{det} integrals: $A_I = \left\langle D_I \middle| \widehat{H}_{\mathsf{TC}} \middle| \Phi \right\rangle$

→ first estimator:

$$a_{I} = \frac{D_{I}}{\Phi \eta^{2}} E_{\text{loc}}^{J} \quad \text{with } E_{\text{loc}}^{J} \equiv \frac{\widehat{H}\left(\Phi e^{J}\right)}{\Phi e^{J}}. \qquad \Rightarrow \langle a_{I} \rangle_{(\eta \Phi)^{2}} \equiv \frac{1}{M} \sum_{i=1}^{M} a_{I}\left(\mathsf{R}_{i}\right) \rightarrow A_{I}$$

→ improved estimator:

$$a_I = \frac{D_I}{\Phi n^2} (E_{\text{loc}}^J - E_{\text{loc}}) + \langle D_I | \widehat{H} | \Phi \rangle$$
 with $E_{\text{loc}} \equiv \frac{\widehat{H}\Phi}{\Phi}$.

→ more improved estimator:

$$a_{I} = \frac{D_{I}}{\Phi n^{2}} (E_{\mathsf{loc}}^{J} - E_{\mathsf{loc}}^{\mathcal{J}}) + \left\langle D_{I} \left| e^{-\mathcal{J}} \widehat{H} e^{\mathcal{J}} \right| \Phi \right\rangle \quad \mathsf{with} \ E_{\mathsf{loc}}^{\mathcal{J}} \equiv \frac{\widehat{H} \left(\Phi e^{\mathcal{J}} \right)}{\Phi e^{\mathcal{J}}}.$$



Statistical fluctuations: illustration for N₂ in cc-pvtz

